

Stochastic explosions in branching processes and non-uniqueness for nonlinear PDE

Radu Dascaliuc

Oregon State University

NSE - quick look

Incompressible **Navier-Stokes Equations:**

$$\left\| \begin{array}{l} \partial_t u + (u \cdot \nabla)u + \nabla p - \Delta u = 0 \\ \nabla \cdot u = 0 \\ u(x, 0) = u_0(x) \text{ (& B.C...)} \end{array} \right.$$

Formally:

$$\left\| \begin{array}{l} u_t + Au + B(u, u) = 0 \\ u(0) = u_0 \end{array} \right.$$

Known: via Approximations/Fixed Point arguments:

- ▶ local existence and uniqueness of *smooth* solutions;
- ▶ global existence and uniqueness for *small initial data*;
- ▶ global existence of *weak* (Leray-Hopf) solutions; *uniqueness?*

Open Problems:

- ▶ Global well-posedness: existence and *uniqueness*...
- ▶ Regularity problem: *Can smooth solutions blow up?*

Scaling and Regularity

$u(x, t)$ – soln. $\Rightarrow u_\lambda(x, t) = \lambda u(\lambda x, \lambda^2 t)$ is also a soln.

What is known globally:

" $\|u(t)\|^2 \leq \|u_0\|^2 < \infty$ " – not enough for regularity.

The most basic regularity criterion:

" $\|A^{\frac{1}{2}} u(t)\|^2 < \infty \Rightarrow$ regularity."

Basic small initial data result:

" $\|A^{\frac{1}{2}} u_0\| \|u_0\| < \epsilon \Rightarrow$ regularity."

Note:

$$\|u_\lambda(t)\|^2 = \frac{1}{\lambda} \|u(\lambda^2 t)\|^2 \text{ – "subcritical"}$$

$$\|A^{\frac{1}{2}} u_\lambda(t)\|^2 = \int_{\mathbb{R}^3} |(-\Delta)^{\frac{1}{2}} \lambda u(\lambda^2 t, \lambda x)|^2 dx = \lambda \|A^{\frac{1}{2}} u(\lambda^2 t)\|^2 \text{ – "supercritical"}$$

$$\|A^{\frac{1}{2}} (u_0)_\lambda\| \|(u_0)_\lambda\| = \|A^{\frac{1}{2}} u_0\| \|u_0\| \text{ – "critical"}$$

NSE - quick look

Incompressible **Navier-Stokes Equations:**

$$\left\| \begin{array}{l} \partial_t u + (u \cdot \nabla)u + \nabla p - \Delta u = 0 \\ \nabla \cdot u = 0 \\ u(x, 0) = u_0(x) \text{ (& B.C...)} \end{array} \right.$$

Formally:

$$\left\| \begin{array}{l} u_t + Au + B(u, u) = 0 \\ u(0) = u_0 \end{array} \right.$$

Known: via Approximations/Fixed Point arguments:

- ▶ local existence and uniqueness of *smooth* solutions;
- ▶ global existence and uniqueness for *small initial data*;
- ▶ global existence of *weak* (Leray-Hopf) solutions; *uniqueness?*

Open Problems:

- ▶ Global well-posedness: existence and *uniqueness*...
- ▶ Regularity problem: *Can smooth solutions blow up?*

Scaling and Regularity

$u(x, t)$ – soln. $\Rightarrow u_\lambda(x, t) = \lambda u(\lambda x, \lambda^2 t)$ is also a soln.

What is known globally:

" $\|u(t)\|^2 \leq \|u_0\|^2 < \infty$ " – not enough for regularity.

The most basic regularity criterion:

" $\|A^{\frac{1}{2}} u(t)\|^2 < \infty \Rightarrow$ regularity."

Basic small initial data result:

" $\|A^{\frac{1}{2}} u_0\| \|u_0\| < \epsilon \Rightarrow$ regularity."

Note:

$$\|u_\lambda(t)\|^2 = \frac{1}{\lambda} \|u(\lambda^2 t)\|^2 \text{ – "subcritical"}$$

$$\|A^{\frac{1}{2}} u_\lambda(t)\|^2 = \int_{\mathbb{R}^3} |(-\Delta)^{\frac{1}{2}} \lambda u(\lambda^2 t, \lambda x)|^2 dx = \lambda \|A^{\frac{1}{2}} u(\lambda^2 t)\|^2 \text{ – "supercritical"}$$

$$\|A^{\frac{1}{2}} (u_0)_\lambda\| \|(u_0)_\lambda\| = \|A^{\frac{1}{2}} u_0\| \|u_0\| \text{ – "critical"}$$

Self-similar solutions:

Definition

u is a self-similar solution if $u(x, t) = u_\lambda(x, t) \forall \lambda > 0$.

In Fourier space:

$$u(x, t) = \lambda u(\lambda x, \lambda^2 t) \Rightarrow \hat{u}(\xi, t) = \frac{1}{\lambda^2} \hat{u}\left(\frac{\xi}{\lambda}, \lambda^2 t\right).$$

Can self-similar Soln's help understand the role of scaling in NSE problem?

Effect of self-similar scaling.

NSE in mild formulation in Fourier space:

$$\widehat{u}(\xi, t) = \widehat{u}_0(\xi) e^{-|\xi|^2 t} + \int_0^t e^{-|\xi|^2 s} \frac{|\xi|}{(2\pi)^{3/2}} \int_{\mathbb{R}^3} \widehat{u}(\eta, t-s) \odot_{e_\xi} \widehat{u}(\xi-\eta, t-s) d\eta ds,$$

with $\xi \cdot \widehat{u}(\xi) = 0$, $\widehat{u}(-\xi) = \overline{\widehat{u}(\xi)}$, $e_\xi = \xi/|\xi|$,

$$v \odot_{e_\xi} w = -i(e_\xi \cdot w) \pi_{e_\xi^\perp} v.$$

In self-similar case, re-scale using $\widehat{u}(\xi, t) = \frac{1}{\lambda^2} \widehat{u}\left(\frac{\xi}{\lambda}, \lambda^2 t\right)$ with $\lambda = |\xi|$:

$$\begin{aligned} \widehat{u}(e_\xi, |\xi|^2 t) &= \widehat{u}_0(e_\xi) e^{-|\xi|^2 t} \\ &+ \frac{1}{(2\pi)^{3/2}} \int_0^\tau \int_{\mathbb{R}^3} e^{-|\xi|^2 s} \widehat{u}(e_\eta, |\eta|^2(t-s)) \odot_\xi \widehat{u}(e_{\xi-\eta}, |\xi-\eta|^2(t-s)) \frac{|\xi|^3 d\eta ds}{|\eta|^2 |\xi-\eta|^2}. \end{aligned}$$

After a change of vars. (with $\tau = |\xi|^2 t$ – "similarity horizon")...

Self-Similar rescaling of NSE

For $\tau = |\xi|^2 t$:

$$\hat{u}(\mathbf{e}_\xi, \tau) = \hat{u}_0(\mathbf{e}_\xi) e^{-\tau} + \frac{1}{(2\pi)^{\frac{3}{2}}} \int_0^\tau \underbrace{e^{-\sigma}}_{\text{Exp}(1)} \int_{\mathbb{R}^3} \hat{u}(\mathbf{e}_\eta, |\eta|^2(\tau - \sigma)) \odot_\xi \hat{u}(\mathbf{e}_{\mathbf{e}_\xi - \eta}, |\mathbf{e}_\xi - \eta|^2(\tau - \sigma)) \underbrace{\frac{d\eta d\sigma}{|\eta|^2 |\mathbf{e}_\xi - \eta|^2}}_{\pi^3 H(\eta|\mathbf{e}_\xi) d\eta d\sigma}$$

In non-ss case the same re-scaling $v(\xi, t) = c \frac{\hat{u}(\xi, t)}{|\xi|^2}$ leads to:

$$v(\xi, t) = v_0(\xi) e^{-|\xi|^2 t} + \int_0^t |\xi|^2 e^{-|\xi|^2 s} \int_{\mathbb{R}^3} v(\eta, t-s) \odot_{\mathbf{e}_\xi} v(\xi - \eta, t-s) H(\eta|\xi) d\eta ds,$$

$$H(\eta|\xi) = \frac{|\xi|}{\pi^3 |\eta|^2 |\xi - \eta|^2}.$$

Probabilistic interpretation.

Abuse notation $\hat{u} \rightsquigarrow v \dots$

$$v(\mathbf{e}_\xi, \tau) = v_0(\mathbf{e}_\xi) e^{-\tau} + \int_0^\tau e^{-\sigma} \int_{\mathbb{R}^3} v(\mathbf{e}_\eta, |\eta|^2(\tau - \sigma)) \odot_{\mathbf{e}_\xi} v(\mathbf{e}_{\mathbf{e}_\xi - \eta}, |\mathbf{e}_\xi - \eta|^2(\tau - \sigma)) \frac{d\eta d\sigma}{\pi^3 |\eta|^2 |\mathbf{e}_\xi - \eta|^2}$$

“Solution” process:

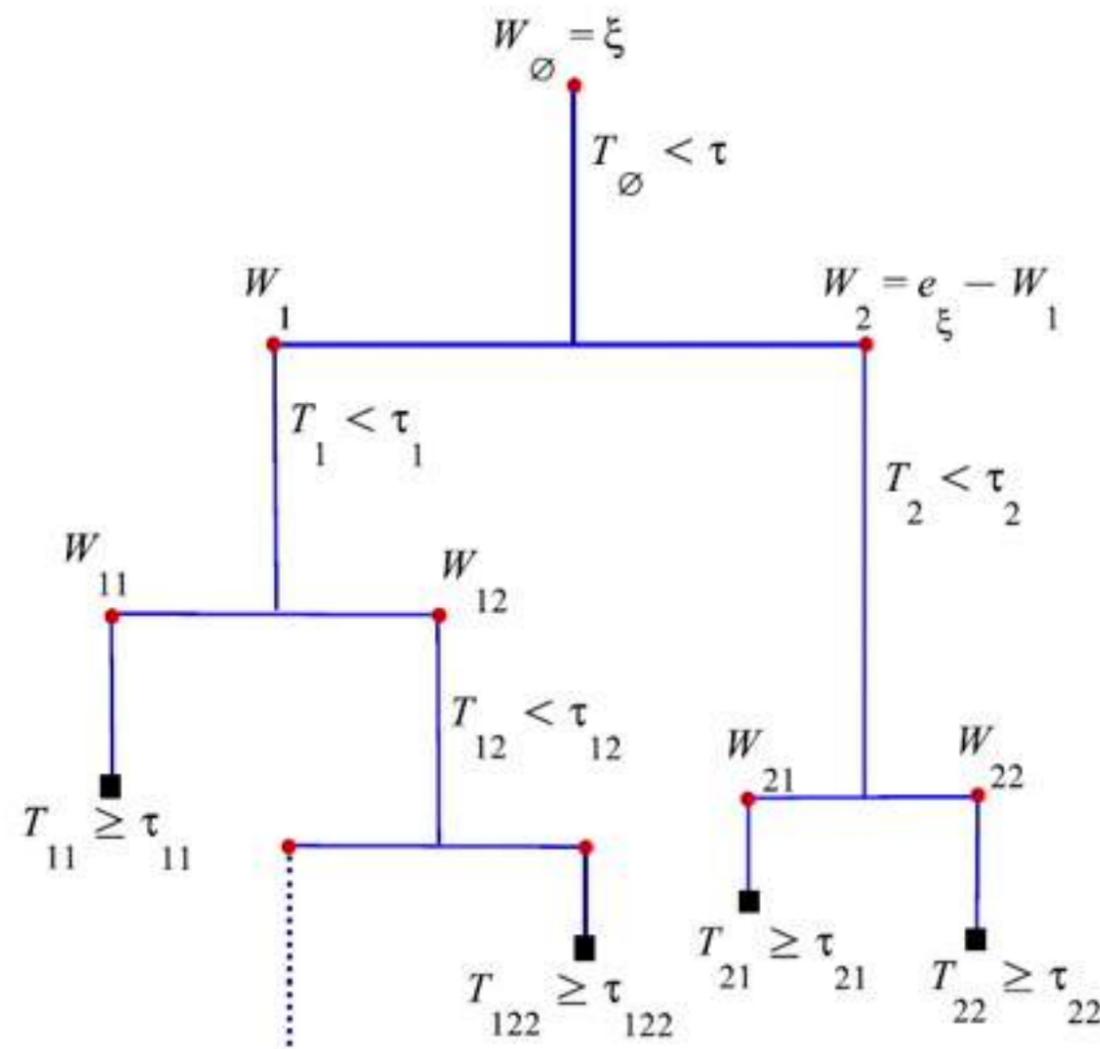
$$X(\mathbf{e}, \tau) = v_0(\mathbf{e}) \mathbf{1}_{T_\emptyset \geq \tau} + X^{(1)}(\mathbf{e}_{W_1}, \tau_1) \odot_{\mathbf{e}} X^{(2)}(\mathbf{e}_{W_2}, \tau_2) \mathbf{1}_{T_\emptyset < \tau},$$

where $T_\emptyset \sim \text{Exp}(1)$, $W_1 \sim H(\cdot | \mathbf{e})$, $W_2 = \mathbf{e} - W_1$, $\tau_j = |W_j|^2(\tau - T_\emptyset)$.

If $\mathbb{E}(|X(\mathbf{e}_\xi, \tau)|) < \infty$, then:

$v(\mathbf{e}_\xi, \tau) = \mathbb{E}(X(\mathbf{e}_\xi, \tau))$ is a solution to ssmNSE.

Self-Similar Cascade.



- ▶ T_b – i.i.d. $\text{Exp}(1)$, $e_b = e_{W_b}$
- ▶ $W_{b1} \sim H(\cdot | e_b)$, $W_{b2} = e_b - W_{b1}$
- ▶ Similarity horizon changes:
 $\tau_{bj} = |W_{bj}|^2 (\tau_b - T_b)$
- ▶ $X(t, e_\xi)$ is the \odot -product of $v_0(e_b)$ s.t.
 $T_b \geq \tau_b$.
- ▶ $v(t, e_\xi) = \mathbb{E}X(t, e_\xi)$ solves ssmNSE if
 $\mathbb{E}|X(t, \xi)| < \infty$
- ▶ **Non-explosion** \Rightarrow uniqueness.

Note: if $\tau_{121} \geq T_{121}$ (the dotted branch in the picture ends):

$$X(e_\emptyset, \tau_\emptyset) = [v_0(e_{11}) \odot_{e_1} (v_0(e_{121}) \odot_{e_{12}} v_0(e_{122}))] \odot_{e_\emptyset} (v_0(e_{21}) \odot_{e_2} v_0(e_{22})).$$

Probabilistic interpretation.

Abuse notation $\hat{u} \rightsquigarrow v \dots$

$$v(\mathbf{e}_\xi, \tau) = v_0(\mathbf{e}_\xi) e^{-\tau} + \int_0^\tau e^{-\sigma} \int_{\mathbb{R}^3} v(\mathbf{e}_\eta, |\eta|^2(\tau - \sigma)) \odot_{\mathbf{e}_\xi} v(\mathbf{e}_{\mathbf{e}_\xi - \eta}, |\mathbf{e}_\xi - \eta|^2(\tau - \sigma)) \frac{d\eta d\sigma}{\pi^3 |\eta|^2 |\mathbf{e}_\xi - \eta|^2}$$

“Solution” process:

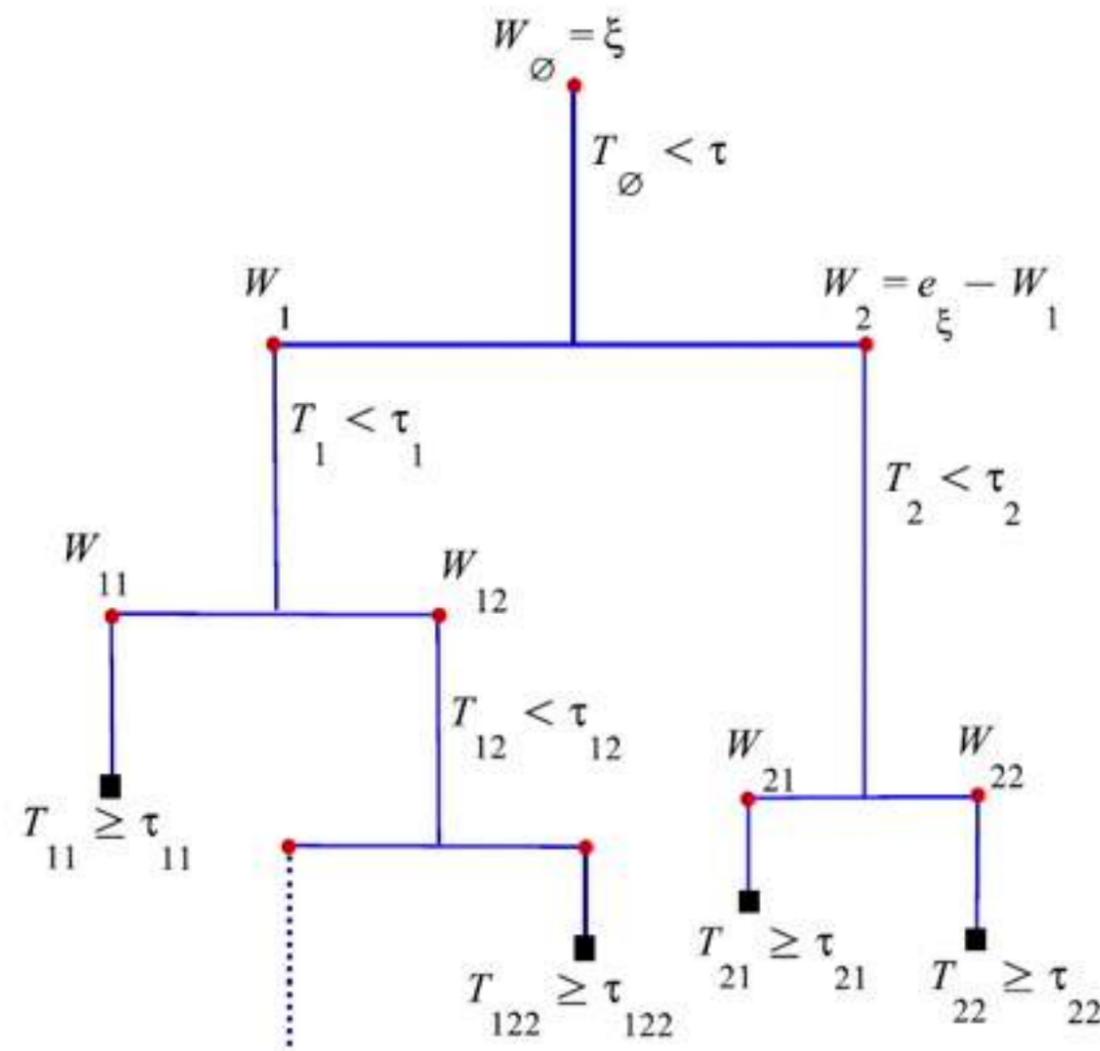
$$X(\mathbf{e}, \tau) = v_0(\mathbf{e}) \mathbf{1}_{T_\emptyset \geq \tau} + X^{(1)}(\mathbf{e}_{W_1}, \tau_1) \odot_{\mathbf{e}} X^{(2)}(\mathbf{e}_{W_2}, \tau_2) \mathbf{1}_{T_\emptyset < \tau},$$

where $T_\emptyset \sim \text{Exp}(1)$, $W_1 \sim H(\cdot | \mathbf{e})$, $W_2 = \mathbf{e} - W_1$, $\tau_j = |W_j|^2(\tau - T_\emptyset)$.

If $\mathbb{E}(|X(\mathbf{e}_\xi, \tau)|) < \infty$, then:

$v(\mathbf{e}_\xi, \tau) = \mathbb{E}(X(\mathbf{e}_\xi, \tau))$ is a solution to ssmNSE.

Self-Similar Cascade.



- ▶ T_b – i.i.d. $\text{Exp}(1)$, $e_b = e_{W_b}$
- ▶ $W_{b1} \sim H(\cdot | e_b)$, $W_{b2} = e_b - W_{b1}$
- ▶ Similarity horizon changes:
 $\tau_{bj} = |W_{bj}|^2 (\tau_b - T_b)$
- ▶ $X(t, e_\xi)$ is the \odot -product of $v_0(e_b)$ s.t.
 $T_b \geq \tau_b$.
- ▶ $v(t, e_\xi) = \mathbb{E}X(t, e_\xi)$ solves ssmNSE if
 $\mathbb{E}|X(t, \xi)| < \infty$
- ▶ **Non-explosion** \Rightarrow uniqueness.

Note: if $\tau_{121} \geq T_{121}$ (the dotted branch in the picture ends):

$$X(e_\emptyset, \tau_\emptyset) = [v_0(e_{11}) \odot_{e_1} (v_0(e_{121}) \odot_{e_{12}} v_0(e_{122}))] \odot_{e_\emptyset} (v_0(e_{21}) \odot_{e_2} v_0(e_{22})).$$

Probabilistic interpretation.

Abuse notation $\hat{u} \rightsquigarrow v \dots$

$$v(\mathbf{e}_\xi, \tau) = v_0(\mathbf{e}_\xi) e^{-\tau} + \int_0^\tau e^{-\sigma} \int_{\mathbb{R}^3} v(\mathbf{e}_\eta, |\eta|^2(\tau - \sigma)) \odot_{\mathbf{e}_\xi} v(\mathbf{e}_{\mathbf{e}_\xi - \eta}, |\mathbf{e}_\xi - \eta|^2(\tau - \sigma)) \frac{d\eta d\sigma}{\pi^3 |\eta|^2 |\mathbf{e}_\xi - \eta|^2}$$

“Solution” process:

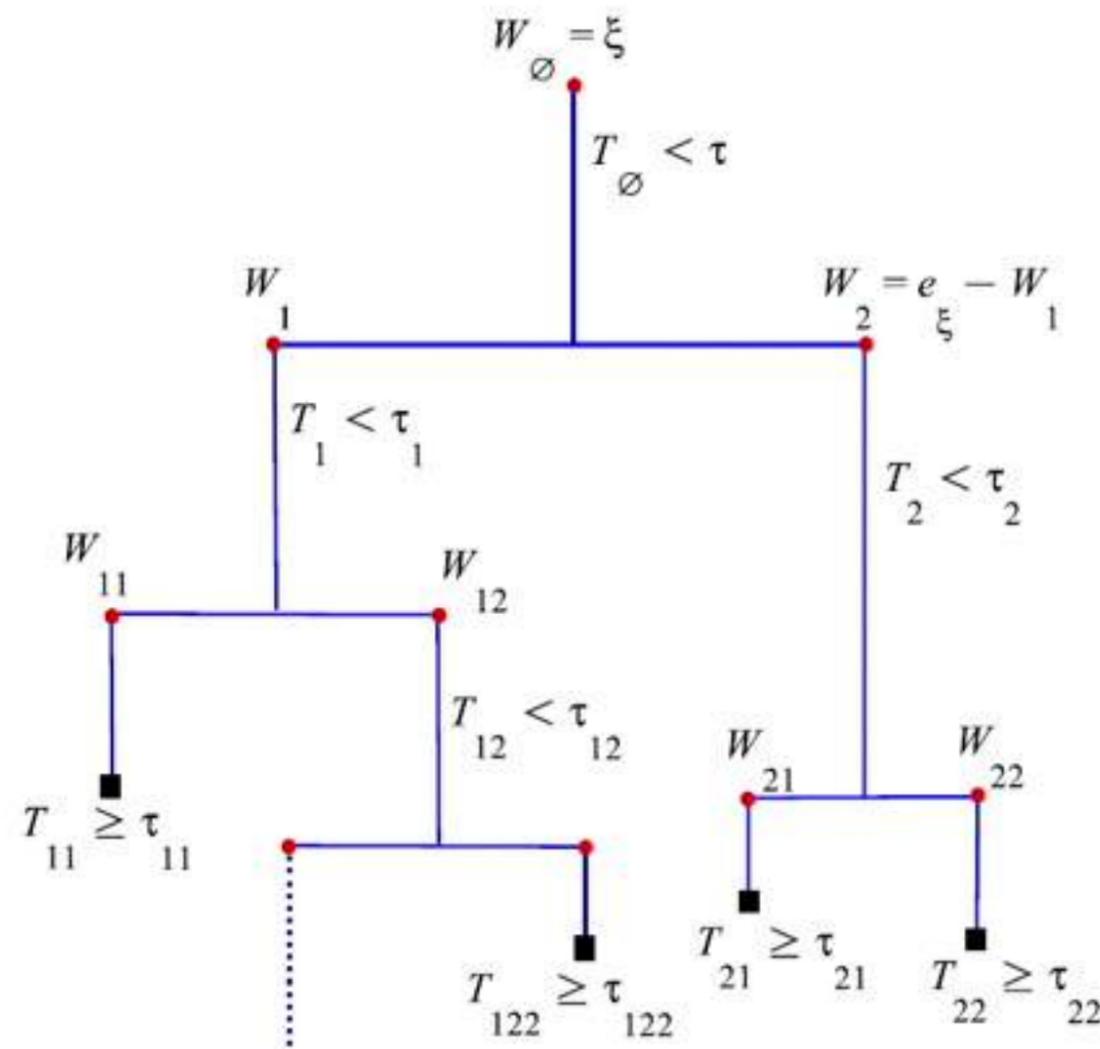
$$X(\mathbf{e}, \tau) = v_0(\mathbf{e}) \mathbf{1}_{T_\emptyset \geq \tau} + X^{(1)}(\mathbf{e}_{W_1}, \tau_1) \odot_{\mathbf{e}} X^{(2)}(\mathbf{e}_{W_2}, \tau_2) \mathbf{1}_{T_\emptyset < \tau},$$

where $T_\emptyset \sim \text{Exp}(1)$, $W_1 \sim H(\cdot | \mathbf{e})$, $W_2 = \mathbf{e} - W_1$, $\tau_j = |W_j|^2(\tau - T_\emptyset)$.

If $\mathbb{E}(|X(\mathbf{e}_\xi, \tau)|) < \infty$, then:

$v(\mathbf{e}_\xi, \tau) = \mathbb{E}(X(\mathbf{e}_\xi, \tau))$ is a solution to ssmNSE.

Self-Similar Cascade.



- ▶ T_b – i.i.d. $\text{Exp}(1)$, $e_b = e_{W_b}$
- ▶ $W_{b1} \sim H(\cdot | e_b)$, $W_{b2} = e_b - W_{b1}$
- ▶ Similarity horizon changes:
 $\tau_{bj} = |W_{bj}|^2 (\tau_b - T_b)$
- ▶ $X(t, e_\xi)$ is the \odot -product of $v_0(e_b)$ s.t.
 $T_b \geq \tau_b$.
- ▶ $v(t, e_\xi) = \mathbb{E}X(t, e_\xi)$ solves ssmNSE if
 $\mathbb{E}|X(t, \xi)| < \infty$
- ▶ **Non-explosion** \Rightarrow uniqueness.

Note: if $\tau_{121} \geq T_{121}$ (the dotted branch in the picture ends):

$$X(e_\emptyset, \tau_\emptyset) = [v_0(e_{11}) \odot_{e_1} (v_0(e_{121}) \odot_{e_{12}} v_0(e_{122}))] \odot_{e_\emptyset} (v_0(e_{21}) \odot_{e_2} v_0(e_{22})).$$

Probabilistic interpretation.

Abuse notation $\hat{u} \rightsquigarrow v \dots$

$$v(\mathbf{e}_\xi, \tau) = v_0(\mathbf{e}_\xi) e^{-\tau} + \int_0^\tau e^{-\sigma} \int_{\mathbb{R}^3} v(\mathbf{e}_\eta, |\eta|^2(\tau - \sigma)) \odot_{\mathbf{e}_\xi} v(\mathbf{e}_{\mathbf{e}_\xi - \eta}, |\mathbf{e}_\xi - \eta|^2(\tau - \sigma)) \frac{d\eta d\sigma}{\pi^3 |\eta|^2 |\mathbf{e}_\xi - \eta|^2}$$

“Solution” process:

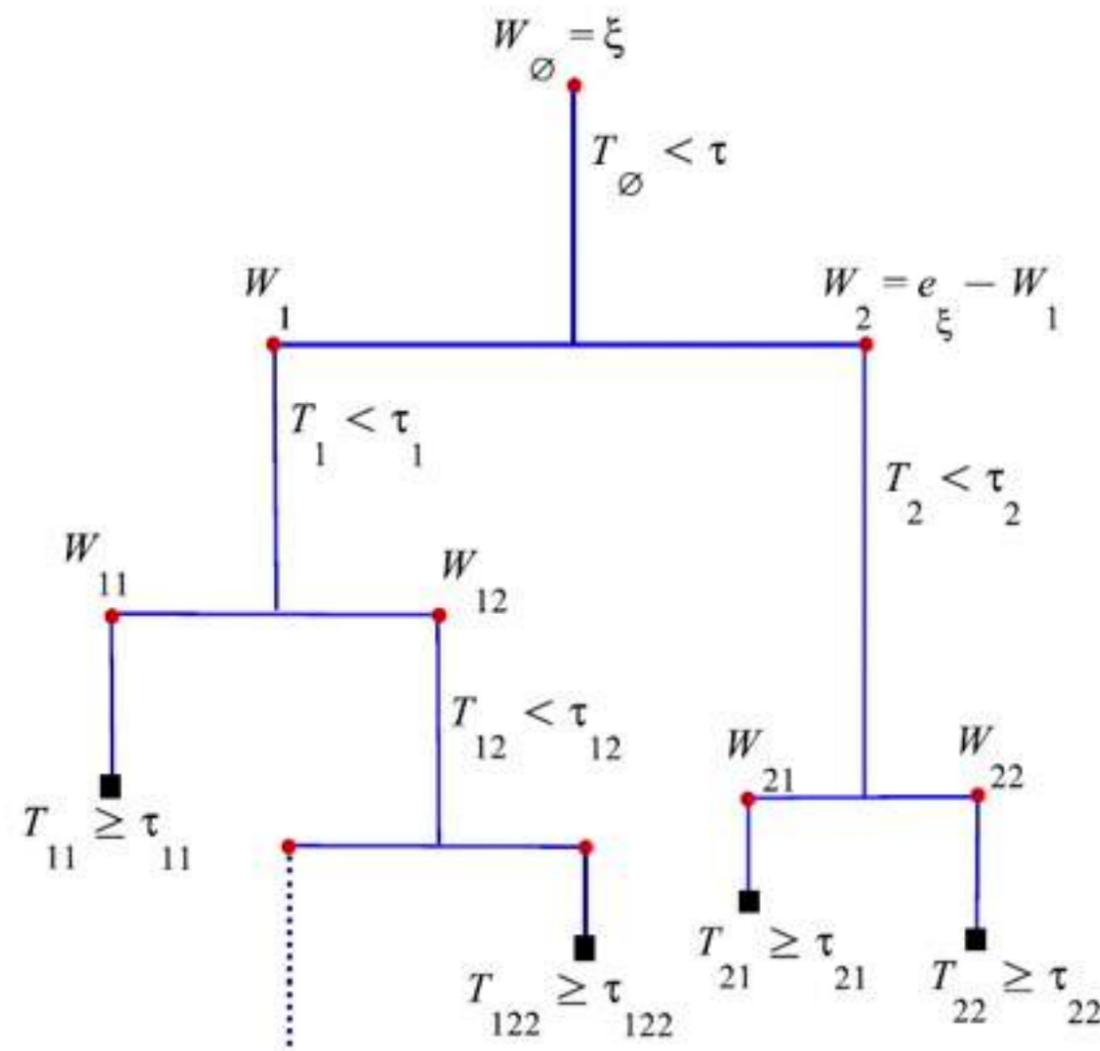
$$X(\mathbf{e}, \tau) = v_0(\mathbf{e}) \mathbf{1}_{T_\emptyset \geq \tau} + X^{(1)}(\mathbf{e}_{W_1}, \tau_1) \odot_{\mathbf{e}} X^{(2)}(\mathbf{e}_{W_2}, \tau_2) \mathbf{1}_{T_\emptyset < \tau},$$

where $T_\emptyset \sim \text{Exp}(1)$, $W_1 \sim H(\cdot | \mathbf{e})$, $W_2 = \mathbf{e} - W_1$, $\tau_j = |W_j|^2(\tau - T_\emptyset)$.

If $\mathbb{E}(|X(\mathbf{e}_\xi, \tau)|) < \infty$, then:

$v(\mathbf{e}_\xi, \tau) = \mathbb{E}(X(\mathbf{e}_\xi, \tau))$ is a solution to ssmNSE.

Self-Similar Cascade.



- ▶ T_b – i.i.d. $\text{Exp}(1)$, $e_b = e_{W_b}$
- ▶ $W_{b1} \sim H(\cdot | e_b)$, $W_{b2} = e_b - W_{b1}$
- ▶ Similarity horizon changes:
 $\tau_{bj} = |W_{bj}|^2 (\tau_b - T_b)$
- ▶ $X(t, e_\xi)$ is the \odot -product of $v_0(e_b)$ s.t.
 $T_b \geq \tau_b$.
- ▶ $v(t, e_\xi) = \mathbb{E}X(t, e_\xi)$ solves ssmNSE if
 $\mathbb{E}|X(t, \xi)| < \infty$
- ▶ **Non-explosion** \Rightarrow uniqueness.

Note: if $\tau_{121} \geq T_{121}$ (the dotted branch in the picture ends):

$$X(e_\emptyset, \tau_\emptyset) = [v_0(e_{11}) \odot_{e_1} (v_0(e_{121}) \odot_{e_{12}} v_0(e_{122}))] \odot_{e_\emptyset} (v_0(e_{21}) \odot_{e_2} v_0(e_{22})).$$

Self-similar explosion

Recall: cascade continues at $b \in \{1, 2\}^n$ if $\tau_b - T_b > 0$. But:

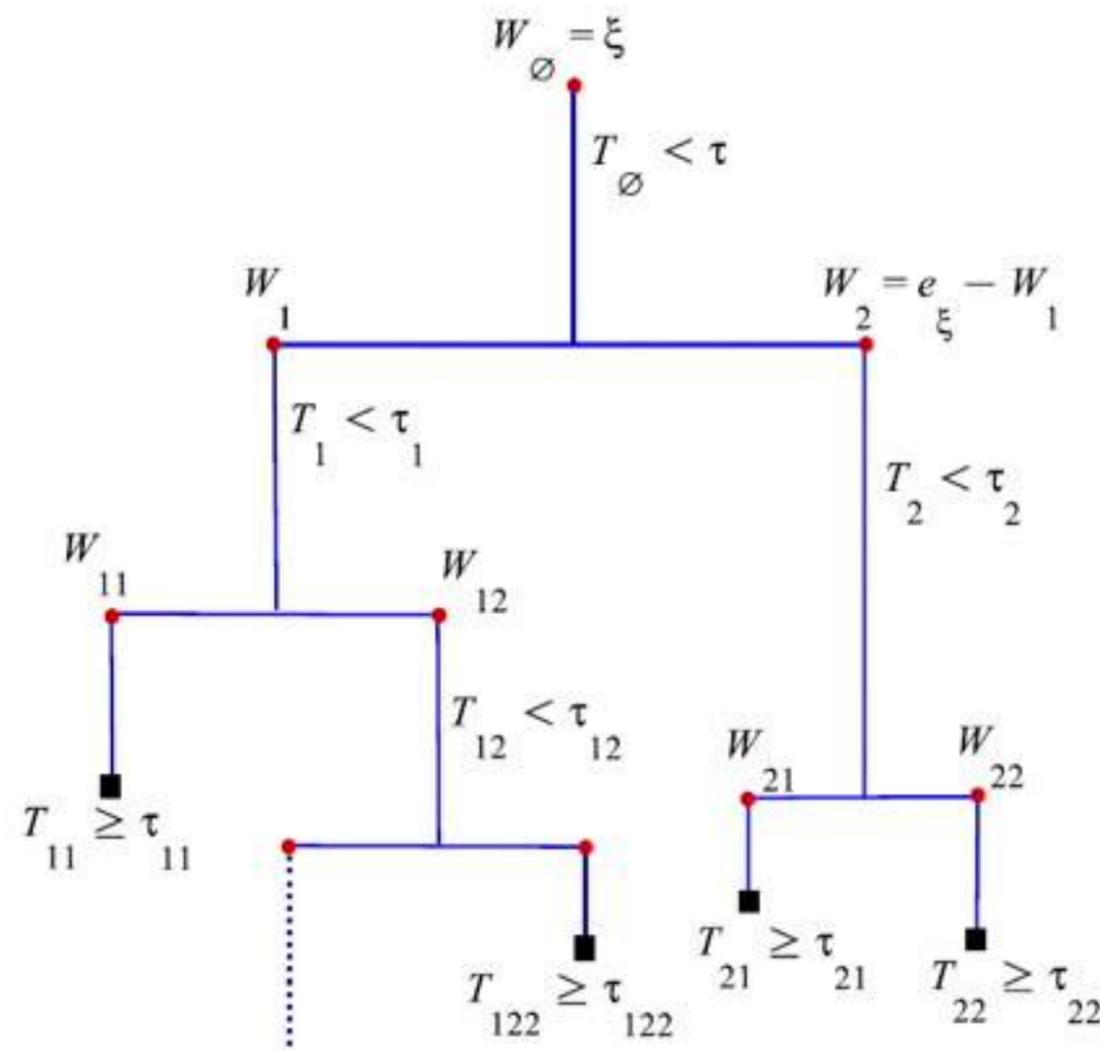
$$\begin{aligned}\tau_b - T_b &= |W_b|^2 (\tau_{b|n-1} - T_{b|n-1}) - T_b = \dots \\ &= \prod_{k=1}^n |W_{b|k}|^2 \left(\tau_\emptyset - \sum_{j=0}^n \frac{T_{b|j}}{\prod_{k=0}^j |W_{b|k}|^2} \right)\end{aligned}$$

\Rightarrow "explosion horizon / shortest path" R.V:

$$S(e_\xi) = \lim_{n \rightarrow \infty} \inf_{|b|=n} \underbrace{\sum_{j=0}^n \frac{T_{b|j}}{\prod_{k=0}^j |W_{b|k}|^2}}_{\theta_b}$$

Explosion event: $E = \{S < \infty\}$.

Self-Similar Cascade.



- ▶ T_b – i.i.d. $\text{Exp}(1)$, $e_b = e_{W_b}$
- ▶ $W_{b1} \sim H(\cdot | e_b)$, $W_{b2} = e_b - W_{b1}$
- ▶ Similarity horizon changes:
 $\tau_{bj} = |W_{bj}|^2 (\tau_b - T_b)$
- ▶ $X(t, e_\xi)$ is the \odot -product of $v_0(e_b)$ s.t.
 $T_b \geq \tau_b$.
- ▶ $v(t, e_\xi) = \mathbb{E}X(t, e_\xi)$ solves ssmNSE if
 $\mathbb{E}|X(t, \xi)| < \infty$
- ▶ **Non-explosion** \Rightarrow uniqueness.

Note: if $\tau_{121} \geq T_{121}$ (the dotted branch in the picture ends):

$$X(e_\emptyset, \tau_\emptyset) = [v_0(e_{11}) \odot_{e_1} (v_0(e_{121}) \odot_{e_{12}} v_0(e_{122}))] \odot_{e_\emptyset} (v_0(e_{21}) \odot_{e_2} v_0(e_{22})).$$

Self-similar explosion

Recall: cascade continues at $b \in \{1, 2\}^n$ if $\tau_b - T_b > 0$. But:

$$\begin{aligned}\tau_b - T_b &= |W_b|^2 (\tau_{b|n-1} - T_{b|n-1}) - T_b = \dots \\ &= \prod_{k=1}^n |W_{b|k}|^2 \left(\tau_\emptyset - \sum_{j=0}^n \frac{T_{b|j}}{\prod_{k=0}^j |W_{b|k}|^2} \right)\end{aligned}$$

\Rightarrow "explosion horizon / shortest path" R.V:

$$S(e_\xi) = \lim_{n \rightarrow \infty} \inf_{|b|=n} \underbrace{\sum_{j=0}^n \frac{T_{b|j}}{\prod_{k=0}^j |W_{b|k}|^2}}_{\theta_b}$$

Explosion event: $E = \{S < \infty\}$.

A simplification: the α -Riccati equation

Recall ssmNSE:

$$v(\mathbf{e}, \tau) = v_0(\mathbf{e}) e^{-\tau} + \int_0^{\tau} e^{-\sigma} \int_{\mathbb{R}^3} v(\mathbf{e}_{\eta_1}, |\eta_1|^2(\tau - \sigma)) \odot_{\mathbf{e}} \underbrace{v(\mathbf{e}_{\eta_2}, |\eta_2|^2(\tau - \sigma))}_{\eta_2 = \mathbf{e} - \eta_1} H(\eta_1 | \mathbf{e}) d\eta_1 d\sigma.$$

Simplifications:

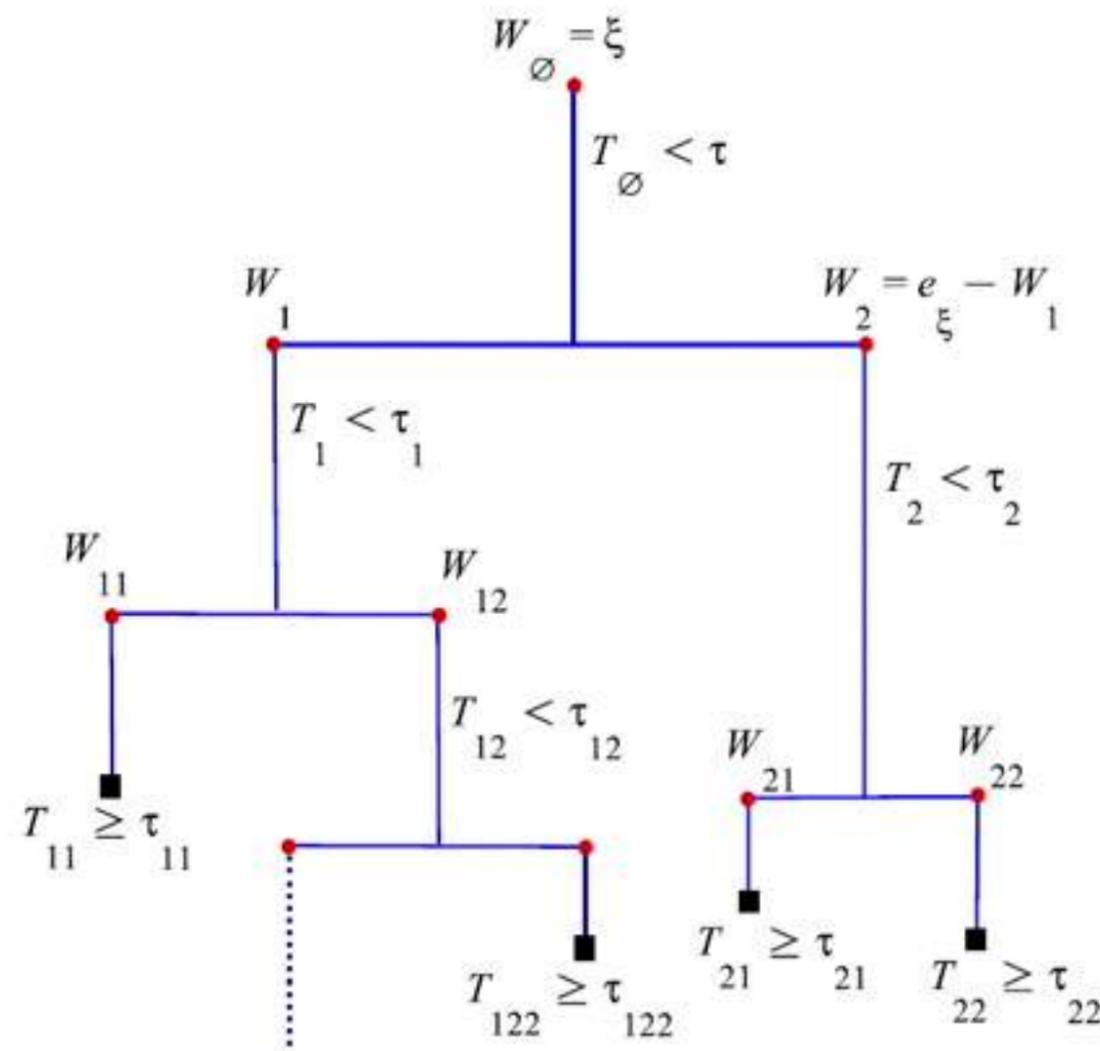
- ▶ $H(\eta_1 | \mathbf{e}) \rightsquigarrow \delta_{\sqrt{\alpha}}$
- ▶ Ignore geometry: $\odot \rightsquigarrow$ product of magnitudes.
- ▶ Ignore the difference between $|\eta_1|$ and $|\eta_2|$.
- ▶ Ignore rotation $|v(\mathbf{e}, \tau)| \rightsquigarrow v(\tau)$.

\Rightarrow α -Riccati equation:

$$v(\tau) = v_0 e^{-\tau} + \int_0^{\tau} e^{-\sigma} v^2(\alpha(\tau - \sigma)) d\sigma.$$

Differential form: $v'(\tau) = -v(\tau) + v^2(\alpha\tau) \dots$

Self-Similar Cascade.



- ▶ T_b – i.i.d. $\text{Exp}(1)$, $e_b = e_{W_b}$
- ▶ $W_{b1} \sim H(\cdot | e_b)$, $W_{b2} = e_b - W_{b1}$
- ▶ Similarity horizon changes:
 $\tau_{bj} = |W_{bj}|^2 (\tau_b - T_b)$
- ▶ $X(t, e_\xi)$ is the \odot -product of $v_0(e_b)$ s.t.
 $T_b \geq \tau_b$.
- ▶ $v(t, e_\xi) = \mathbb{E}X(t, e_\xi)$ solves ssmNSE if
 $\mathbb{E}|X(t, \xi)| < \infty$
- ▶ **Non-explosion** \Rightarrow uniqueness.

Note: if $\tau_{121} \geq T_{121}$ (the dotted branch in the picture ends):

$$X(e_\emptyset, \tau_\emptyset) = [v_0(e_{11}) \odot_{e_1} (v_0(e_{121}) \odot_{e_{12}} v_0(e_{122}))] \odot_{e_\emptyset} (v_0(e_{21}) \odot_{e_2} v_0(e_{22})).$$

Self-similar explosion

Recall: cascade continues at $b \in \{1, 2\}^n$ if $\tau_b - T_b > 0$. But:

$$\begin{aligned}\tau_b - T_b &= |W_b|^2 (\tau_{b|n-1} - T_{b|n-1}) - T_b = \dots \\ &= \prod_{k=1}^n |W_{b|k}|^2 \left(\tau_\emptyset - \sum_{j=0}^n \frac{T_{b|j}}{\prod_{k=0}^j |W_{b|k}|^2} \right)\end{aligned}$$

\Rightarrow "explosion horizon / shortest path" R.V:

$$S(e_\xi) = \lim_{n \rightarrow \infty} \inf_{|b|=n} \underbrace{\sum_{j=0}^n \frac{T_{b|j}}{\prod_{k=0}^j |W_{b|k}|^2}}_{\theta_b}$$

Explosion event: $E = \{S < \infty\}$.

A simplification: the α -Riccati equation

Recall ssmNSE:

$$v(\mathbf{e}, \tau) = v_0(\mathbf{e}) e^{-\tau} + \int_0^\tau e^{-\sigma} \int_{\mathbb{R}^3} v(\mathbf{e}_{\eta_1}, |\eta_1|^2(\tau - \sigma)) \odot_{\mathbf{e}} \underbrace{v(\mathbf{e}_{\eta_2}, |\eta_2|^2(\tau - \sigma))}_{\eta_2 = \mathbf{e} - \eta_1} H(\eta_1 | \mathbf{e}) d\eta_1 d\sigma.$$

Simplifications:

- ▶ $H(\eta_1 | \mathbf{e}) \rightsquigarrow \delta_{\sqrt{\alpha}}$
- ▶ Ignore geometry: $\odot \rightsquigarrow$ product of magnitudes.
- ▶ Ignore the difference between $|\eta_1|$ and $|\eta_2|$.
- ▶ Ignore rotation $|v(\mathbf{e}, \tau)| \rightsquigarrow v(\tau)$.

\Rightarrow α -Riccati equation:

$$v(\tau) = v_0 e^{-\tau} + \int_0^\tau e^{-\sigma} v^2(\alpha(\tau - \sigma)) d\sigma.$$

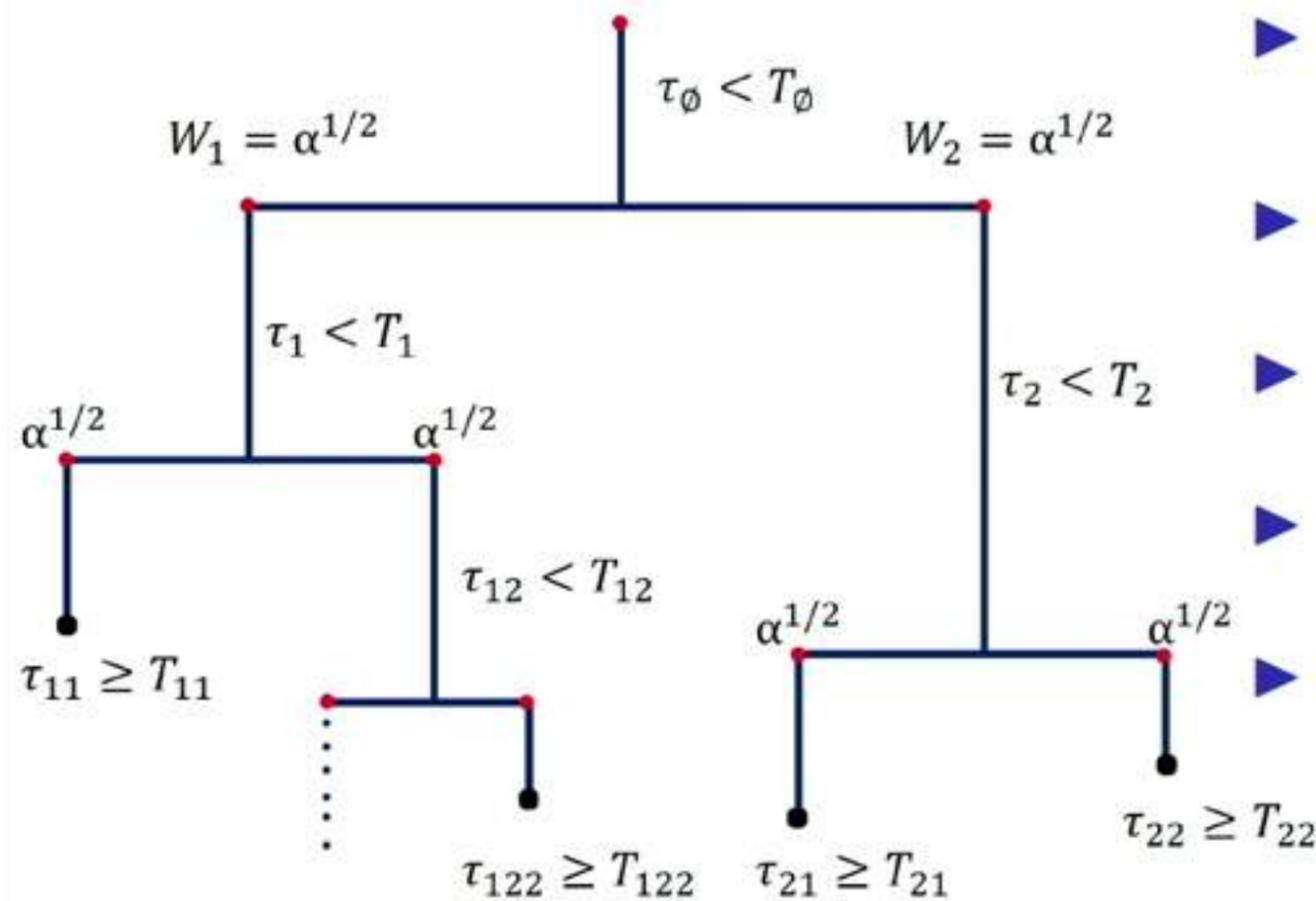
Differential form: $v'(\tau) = -v(\tau) + v^2(\alpha\tau) \dots$

Cascade set-up for α -Riccati:

$$v(\tau) = v_0 e^{-\tau} + \int_0^\tau e^{-\sigma} v^2(\alpha(\tau - \sigma)) d\sigma.$$

Solution Process = Process with *inherited property*:

$$X(\tau) = X(0) \mathbf{1}_{T_\emptyset \geq \tau} + X^{(1)}(\alpha\tau_1) X^{(2)}(\alpha\tau_2) \mathbf{1}_{T_\emptyset < \tau},$$



- ▶ $\mathbb{E}(|X(\tau)|) < \infty \Rightarrow v(\tau) = \mathbb{E}(X(\tau))$ a solution, $v_0 = \mathbb{E}(X(0))$.
- ▶ T_b – i.i.d. $\text{Exp}(1)$.
- ▶ $\tau_{bj} = \alpha(\tau_b - T_b)$, $\tau_\emptyset = \tau$.
- ▶ $T_b \geq \tau_b$ – process at branch b stops.
- ▶ **Non-explosion** \Rightarrow uniqueness.

Note: if $\tau_{121} \geq T_{121}$: $X(\tau_\emptyset) = v_0^5$.

Some inherited events

Define:

$$\text{Shortest path: } S = \lim_{n \rightarrow \infty} \min_{|b|=n} \sum_{k=0}^n \frac{T_{b|k}}{\alpha^k},$$

$$\text{Longest path: } L = \lim_{n \rightarrow \infty} \max_{|b|=n} \sum_{k=0}^n \frac{T_{b|k}}{\alpha^k},$$

explosion events: $\{S < \infty\}$ – explosion; $\{L < \infty\}$ – *hyper-explosion*.

Note:

- ▶ no-explosion and hyper-explosion are inherited events.
- ▶ $X(t) = \mathbb{1}_{S \geq \tau}$, $X(t) = \mathbb{1}_{L < \tau}$ – processes with inherited property.

Conclusion:

$v = \mathbb{P}(S = \infty)$, $v = \mathbb{P}(L < \infty)$, $v(\tau) = \mathbb{P}(S \geq \tau)$, $v(\tau) = \mathbb{P}(L < \tau)$
solve α -Riccati.

The case $\alpha \leq 1$.

Theorem

When $\alpha \in [0, 1]$ the cascade is non-exploding: $\mathbb{P}(S = \infty) = 1$.

Proof.

Compare to Yule process. □

Consequence: $X(t) = v_0^{N(\tau)}$, $N(\tau)$ —number of branches;

Theorem

Let $\alpha < 1$. $\forall v_0 > 0 \exists!$ solution $v(t) = \mathbb{E}(v_0^{N(\tau)}) < \infty$:

$$\lim_{n \rightarrow \infty} \mathbb{E}(v_0^{N(\tau)}) = \begin{cases} 0, & 0 \leq v_0 < 1, \\ 1, & v_0 = 1, \\ \infty, & v_0 > 1 \end{cases} .$$

Proof.

Estimate $\mathbb{P}(N(\tau) = n)$. □

Special Cases.

- ▶ $\alpha = 0$: equation is $v'(\tau) = -v(\tau) + v_0^2$.

$$\mathbb{P}(N(\tau) = 1) = e^{-\tau}, \quad \mathbb{P}(N(\tau) = 2) = 1 - e^{-\tau},$$

$$v(\tau) = \mathbb{E}(v_0^{N(\tau)}) = v_0 e^{-\tau} + v_0^2 (1 - e^{-\tau}).$$

- ▶ $\alpha = 1$: equation is $v'(\tau) = -v(\tau) + v^2(\tau)$.

$$\mathbb{P}(N(\tau) = n) = e^{-\tau} (1 - e^{-\tau})^{n-1}, \quad n \in \mathbb{N}.$$

$$v(\tau) = \mathbb{E}(v_0^{N(\tau)}) = \sum_{n=1}^{\infty} v_0^n e^{-\tau} (1 - e^{-\tau})^{n-1} = \frac{v_0 e^{-\tau}}{1 - v_0(1 - e^{-\tau})}.$$

- ▶ $\alpha = 1/2$: equation is $v'(\tau) = -v(\tau) + v^2(\tau/2)$.

$$\mathbb{P}(N(\tau) = n) = e^{-\tau} \tau^{n-1} / (n-1)!, \quad n \in \mathbb{N}.$$

$$v(\tau) = \mathbb{E}(v_0^{N(\tau)}) = \sum_{n=1}^{\infty} v_0^n \frac{\tau^{n-1}}{(n-1)!} e^{-\tau} = v_0 e^{v_0 \tau - 1}.$$

The case $\alpha > 1$: explosion.

Theorem (Maximal-path Explosion Theorem)

Suppose $\exists C < 1$ s.t. in the self-similar cascade for any b , $Z_b = \max\{|W_{b1}|, |W_{b2}|\}$ satisfies $\mathbb{E}(Z_b^{-2}) \leq C$. Then the cascade is exploding: $\mathbb{P}(S < \infty) = 1$.

...In the α -Riccati case $|W_b| = \alpha \forall b$, so the cascade is **exploding** when $\alpha > 1$.

In fact:

Theorem

When $\alpha > 1$ the cascade for α -Riccati is **hyper-exploding**:
 $\mathbb{P}(L < \infty) = 1$.

Question: How to build process $X(t)$ in the presence exploding branches?

$\alpha > 1$. Iterative Process.

Start with an $X_0(\tau) \geq 0$, and set

$$X_n(\tau) = \begin{cases} v_0, & T_\emptyset \geq \tau, \\ X_{n-1}^{(1)}(\alpha(\tau - T_\emptyset)) X_{n-1}^{(2)}(\alpha(\tau - T_\emptyset)), & T_\emptyset < \tau \end{cases}, \quad n \in \mathbb{N},$$

If $y_n(\tau) = \mathbb{E}(X_n(\tau)) < \infty$:

$$y_n(\tau) = u_0 e^{-\tau} + \int_0^\tau e^{-\sigma} y_{n-1}^2(\alpha(\tau - \sigma)) d\sigma,$$

Question: For which X_0 we can pass to the limit?

$\alpha > 1$. Minimal Solution: Existence.

Choose $X_0 = 0$. Then $X_n(\tau) \rightarrow \underline{X}(\tau) = v_0^{N(t)} \mathbf{1}_{S>t}$.

Note: $v_0 = 1$, $\underline{X}(\tau) = \mathbf{1}_{S>t} \rightarrow 0$, $\tau \rightarrow \infty$; $v_0 = 0$, $\underline{X} = 0$ a.s.

Theorem

- ▶ If $\underline{v}(\tau) = \mathbb{E}(\underline{X}(\tau)) < \infty$, then $\underline{v}(\tau)$ is a solution;
- ▶ if $v(\tau)$ is another solution for the same initial data, then $\underline{v}(\tau) < v(\tau)$.

Theorem

Let $\alpha > 1$.

- ▶ If $v_0 \leq \max\{1, (2\alpha - 1)/4\}$, then $\underline{v}(\tau) = \mathbb{E}(\underline{X}(\tau))$ is a (global) solution.
- ▶ If $v_0 > 2\alpha - 1$ then local solutions blow up in finite time.

$\alpha > 1$: Uniqueness Problem for $0 \leq v_0 \leq 1$.

Theorem

Let $0 \leq v_0 < 1$. Then $\underline{v}(\tau)$ is the *unique solution* in the class $\|v(\tau)\|_\infty < 1$.

Proof.

Let $v(\tau)$ be a solution. Initialize iterations with $X_0 = v(\tau)$ and compare to \underline{X} ...

□

“Upper Solution”: use $X_0 = 1$: $X_n(\tau) \rightarrow \bar{X}(\tau) = v_0^{N(\tau)}$. Set $\bar{v}(\tau) = \mathbb{E}(\bar{X}(\tau))$.

Note: $v_0 = 0$, $\bar{X}(\tau) = \mathbb{1}_{L \leq t} \rightarrow 1$, $\tau \rightarrow \infty$; $v_0 = 1$, $\bar{X} = 1$ a.s.

Theorem

Let $\alpha > 1$ and $v_0 \in [0, 1]$. There are *at least two solutions*:

- ▶ $\underline{v}(\tau) \rightarrow 0$ as $\tau \rightarrow \infty$,
- ▶ $\bar{v}(\tau) \rightarrow 1$ as $\tau \rightarrow \infty$.

Moreover, if $v(\tau) \in [0, 1]$, then $\underline{v}(\tau) \leq v(\tau) \leq \bar{v}(\tau)$.

$\alpha > 1$. Iterative Process.

Start with an $X_0(\tau) \geq 0$, and set

$$X_n(\tau) = \begin{cases} v_0, & T_\emptyset \geq \tau, \\ X_{n-1}^{(1)}(\alpha(\tau - T_\emptyset)) X_{n-1}^{(2)}(\alpha(\tau - T_\emptyset)), & T_\emptyset < \tau \end{cases}, \quad n \in \mathbb{N},$$

If $y_n(\tau) = \mathbb{E}(X_n(\tau)) < \infty$:

$$y_n(\tau) = u_0 e^{-\tau} + \int_0^\tau e^{-\sigma} y_{n-1}^2(\alpha(\tau - \sigma)) d\sigma,$$

Question: For which X_0 we can pass to the limit?

$\alpha > 1$: Uniqueness Problem for $0 \leq v_0 \leq 1$.

Theorem

Let $0 \leq v_0 < 1$. Then $\underline{v}(\tau)$ is the *unique solution* in the class $\|v(\tau)\|_\infty < 1$.

Proof.

Let $v(\tau)$ be a solution. Initialize iterations with $X_0 = v(\tau)$ and compare to \underline{X} ... □

“Upper Solution”: use $X_0 = 1$: $X_n(\tau) \rightarrow \bar{X}(\tau) = v_0^{N(\tau)}$. Set $\bar{v}(\tau) = \mathbb{E}(\bar{X}(\tau))$.

Note: $v_0 = 0$, $\bar{X}(\tau) = \mathbb{1}_{L \leq t} \rightarrow 1$, $\tau \rightarrow \infty$; $v_0 = 1$, $\bar{X} = 1$ a.s.

Theorem

Let $\alpha > 1$ and $v_0 \in [0, 1]$. There are *at least two solutions*:

- ▶ $\underline{v}(\tau) \rightarrow 0$ as $\tau \rightarrow \infty$,
- ▶ $\bar{v}(\tau) \rightarrow 1$ as $\tau \rightarrow \infty$.

Moreover, if $v(\tau) \in [0, 1]$, then $\underline{v}(\tau) \leq v(\tau) \leq \bar{v}(\tau)$.

$\alpha > 1$. Non-uniqueness for big initial data.

Theorem

Let $\alpha > 5/2$ and $1 < u_0 \leq \frac{2\alpha-1}{4} - \phi_\infty$, ($\phi_\infty(\alpha) < 3/4$). Then

$$\underline{v}(\tau) < \bar{v}(\tau) \leq \underline{w} + (1 - e^{-\tau}),$$

where \underline{w} is the minimal solution corresponding to initial data $w_0 \in [v_0 + \phi_\infty, (2\alpha - 1)/4]$.

Also can be shown: local in time solutions are not unique.

A Random Initialization.

$$X_0(t) = \begin{cases} 0, & T_\emptyset \geq \tau, \\ G(\tau - T_\emptyset), & T_\emptyset < \tau, \end{cases}$$

Denote

$$F(t) := \mathbb{E}(X_0(\tau)) = \int_0^\tau e^{-\sigma} G(\tau - \sigma) d\sigma \quad \text{or,} \quad \begin{cases} F'(\tau) + F(\tau) = G(\tau) \\ F(0) = 0. \end{cases}$$

$$\text{If } L_n := \max_{|b|=n} \underbrace{\sum_{k=0}^n \frac{T_{b|k}}{\alpha^k}}_{\theta_b} < \tau \text{ then } X_n(\tau) = \prod_{|\tau|=n} G(\alpha^n(\tau - \theta_b)).$$

and so

$$\mathbb{E}(X_{n+1}(\tau) \mathbf{1}_{L_{n+1} < \tau} | T_b, |b| \leq n) = \prod_{|b|=n} F^2(\alpha^{n+1}(\tau - \theta_b)).$$

Note: if $F^2(\alpha\tau) \leq G(\tau)$, then $X_n(\tau) \mathbf{1}_{L_n < \tau}$ is a super-martingale.

Athreya's Solution.

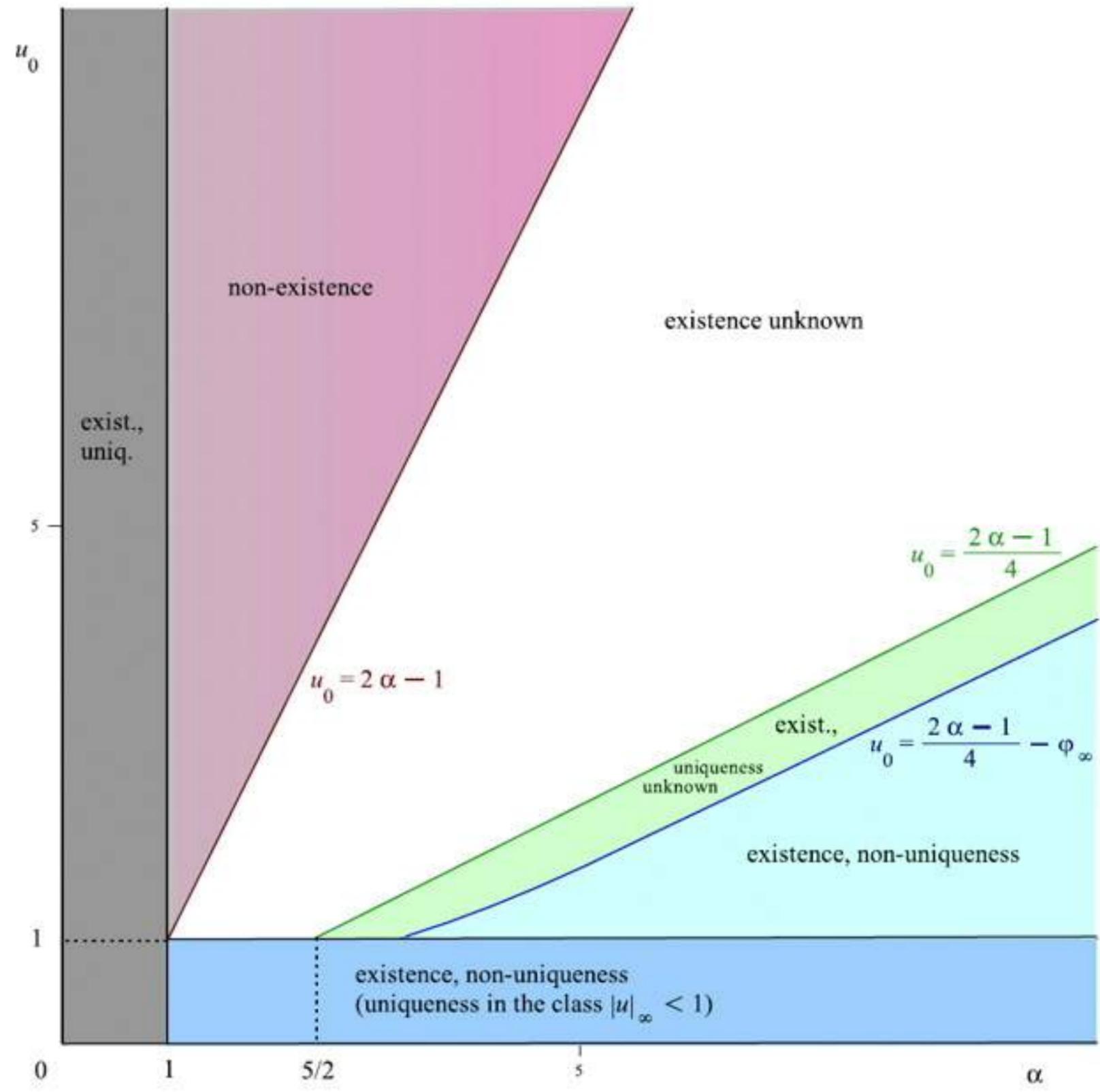
For $\alpha > 1$ and $v_0 = 0$, Choose

$$F(\tau) = e^{-\tau^{-b}}, \quad b = \frac{\ln 2}{\ln \alpha}.$$

Then $X_n(\tau) = X_n(\tau)\mathbf{1}_{L_n < \tau}$ is an *uniformly integrable* super-martingale.

Conclusion: $\exists X(t) = \lim_n X_n(t)$ and $v_A(t) = \mathbb{E}(X(t))$ is yet another solution to α -Riccati.

Existence & Uniqueness regions for α -Riccati.



Back to NSE: explosion?

Some properties of $H(\eta|\mathbf{e}_\xi) = \frac{1}{\pi^3} \frac{1}{|\eta|^2 |\mathbf{e}_\xi - \eta|}$:

- ▶ Rotational invariance;
- ▶ $\forall b, \{|W_{b|j}|\}_j - \text{iid } \frac{2}{\pi^2} \frac{1}{r} \ln \left| \frac{1+r}{1-r} \right|$.
- ▶ In polar coordinates at each branching b :
 - ▶ azimuth angle for W_{b1} is $\text{Unif}(0, 2\pi)$;
 - ▶ longitudinal angles for W_{b1} and W_{b2} are jointly $\text{Unif}\{0 \leq \phi_1, \phi_2 \leq \pi, \phi_1 + \phi_2 \leq \pi\}$.

Consequence:

$$\max\{|W_{b1}|, |W_{b2}|\} \text{ is distributed } z(r) = \frac{4}{\pi} \frac{1}{r} \ln \left| \frac{r}{r-1} \right|.$$

Note:

$$\text{if } Z \sim z(r), \text{ then } \mathbb{E}(Z^{-2}) = \frac{2}{\pi} < 1.$$

Thus, by the Maximal-Path Explosion Theorem ssmNSE cascade is **exploding!**

Maximal-Path Explosion

Recall the Maximal-Path Explosion Theorem:

Theorem

Suppose $\exists C < 1$ s.t. in the self-similar cascade for any b , $Z_b = \max\{|W_{b1}|, |W_{b2}|\}$ satisfies $\mathbb{E}(Z_b^{-2}) \leq C$. Then the cascade is exploding $\mathbb{P}(S < \infty) = 1$.

Idea of Proof:

Recall: Explosion event $\{S = \infty\}$, where $S = \lim_{n \rightarrow \infty} \min_{|b|=n} \sum_{j=0}^n \frac{T_{b|j}}{\prod_{k=0}^j |W_{b|k}|^2}$

Consider “maximal-descendant” path $b_m \in \{1, 2\}^{\mathbb{N}}$ that at each junction $b = b_m|j$ follows the path of maximal $|W_{bi}|$, $i = 1, 2$. Note:

$$S \leq \sum_{j=0}^{\infty} \frac{T_{b_m|j}}{\prod_{k=0}^j |W_{b_m|k}|^2} \quad \text{a.s.}$$

What could be proven for NSE?

Existence and uniqueness for v_0 via **minimal solution**:

Theorem

Suppose $\|v_0(\mathbf{e}_\xi)\|_\infty < 1$. Then there exists a unique ssmNSE solution in the class $\|v(\mathbf{e}_\xi, \tau)\|_\infty < 1$.

(In fact existence can be proven for $\|v_0(\mathbf{e}_\xi)\|_\infty \leq C$ for a $C > 1$.)

Conclusions/Challenges

- ▶ Can iterative procedure be used to prove non-uniqueness for NSE? Lack of hyper-explosion is a challenge.
- ▶ Can we prove existence for NSE for arbitrary large initial data?
Good: NSE nonlinearity – \odot -product along a tree – contains repeated projection of the same vector on random planes, which should deplete non-linearity.
Bad: A large tree generates long \odot -products: enough v_0 may survive the depletion and make X large.
- ▶ Can this explosion approach be connected to Kolmogorov backwards equation... Develop a general theory of explosion-nonuniqueness for nonlinear PDE?
- ▶ Can cascade structure be employed to describe energy transfer between scales with consequences in turbulence/regularity?